# Study On Non Linear Equation For Boson Stars 

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#### Abstract

Meanwhile the equation cannot be solved by arithmetic undoing and requires algebraic operations to be performed to simplify the equation to give a solution. This phenomenon is called 'the didactic cut'. It relates to the observation that many students see the 'equals' sign as an operation, arising out of experience in arithmetic where an equation of the form is seen as a dynamic operation to perform the calculation, 'three plus four makes 7', so that an equation such as is seen as an operation which may possibly be solved by arithmetic 'undoing' rather than requiring algebraic manipulation (Kieran, 1981). [Kumar, R. and Yadav, A. Study On Non Linear Equation For Boson Stars. Academ Arena 2020;12(9):13-15]. ISSN 1553-992X (print); ISSN 2158-771X (online). http://www.sciencepub.net/academia. 2. doi:10.7537/marsaaj120920.02.


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## Introduction:

The Historical Roots of Elementary Mathematics (Bunt, Jones, \& Bedient, 1976) is very similar in style and information to Math through the Ages. Both books present information in short chapters specific to a main idea (e.g. Greek numeration systems). In addition, both books cover a wide range of topics that are broken down by date. However, The Historical Roots of Elementary Mathematics does not delve into the stories describing the people behind the discoveries. The four volume collection. The World of Mathematics (Newman, 1956) consists of individual articles compiled together in an effort to convey the "...diversity, the utility and the beauty of mathematics" (Newman, iii). Newman attempted to show the richness and range of mathematics. This collection spans ideas from the Rhind Papyrus to the "Statistics of Deadly Quarrels" (Newman, 1956).

The World of Mathematics presents an amazingly broad view of the many applications of mathematics to the sciences. An Introduction to the History of Math (Eves, 1956) covers the same topics as several of the other books, in much the same manner. It traces the development of mathematics from numeration systems through to the development of calculus. It includes specific information of the individuals that developed many of the critical ideas in the history of mathematics. Boyer's (1968) A History of Mathematics is almost entirely about Greek mathematics. It covers ancient Greek mathematics to a degree that none of the other mentioned texts do. Perhaps one of the most valuable tools for a secondary teacher available is Historical Topics for the Mathematics Classroom (National

Council for Teachers of Mathematics, 1989). This text consists of a series of "capsules" (short chapters). Each capsule gives a brief historical overview of a particular topic (e.g. Napier's Rods). The capsules are grouped by general topic (algebra, geometry, trigonometry, etc.). Specifically, this text provides a historical context to graphical approaches to equation solving. In addition, it provides a concise overview of the methods employed to solve quadratics and cubics.

Various researchers (Vaiyavutjamai \& Clements, 2006) have illustrated that very little attention has been paid to quadratic equations in mathematics education literature, and there is scarce research regarding the teaching and learning of quadratic equations. A limited number of research studies focusing on quadratic equations have documented the techniques students engage in while solving quadratic equations (Bossé \& Nandakumar, 2005), geometric approaches used by students for solving quadratic equations (Allaire \& Bradley, 2001), students' understanding of and difficulties with solving quadratic equations (Kotsopoulos, 2007; Lima, 2008; Tall, Lima, \& Healy, 2014; Vaiyavutjamai, Ellerton, \& Clements, 2005; Zakaria \& Maat, 2010), the teaching and learning of quadratic equations in classrooms (Olteanu \& Holmqvist, 2012; Vaiyavutjamai \& Clements, 2006), comparing how quadratic equations are handled in mathematics textbooks in different countries (Saglam \& Alacaci, 2012), and the application of the history of quadratic equations in teacher preparation programs to highlight prospective teachers' knowledge (Clark,
2012). In general, for most students, quadratic equations create challenges in various ways such as difficulties in algebraic procedures, (particularly in factoring quadratic equations), and an inability to apply meaning to the quadratics. Kotsopoulos (2007) suggests that recalling main multiplication facts directly influences a student's ability while engaged in factoring quadratics. Furthermore, since solving the quadratic equations by factorization requires students to find factors rapidly, factoring simple quadratics becomes quite a challenge, while non-simple ones (i.e., $\mathrm{ax} 2+\mathrm{bx}+\mathrm{c}$ where $\mathrm{a}^{\wedge} 1$ ) become harder still. Factoring quadratics can be considerably complicated when the leading coefficient or the constant term has many pairs of factors (Bossé \& Nandakumar, 2005).

## Overview of Methodology

The nonlinear dispersive condition was thoroughly determined as of late in for a boson star, which alludes to a quantum mechanical arrangement of N bosons with relativistic scattering connecting
through a gravitational alluring or appalling Coulomb potential. Indeed, by beginning from the N -body relativistic Schrodinger condition (supplanting $-\Delta / 2$ in the Schrodinger equation (2.1) to $\sqrt{-\Delta+m^{2}}$ ) what's more, picking the underlying wave capacity to portray a condensate where N bosons are all in a similar one-molecule state, in the mean-field limit N $\rightarrow \infty$, one can demonstrate that the time advancement of the one-molecule thickness is represented by the nonlinear relativistic Hartree condition (under a legitimate non-dimensionalization).

## Non Linear Equation For Boson Stars

It is anything but difficult to demonstrate that the condition (1.8) concedes at any rate two significant saved amounts for example the mass of the framework $\mathrm{N}(\psi(\cdot, \mathrm{t})):=\mathrm{k} \psi(\cdot, \mathrm{t}) \mathrm{k} 2=\mathrm{Z} 3|\psi(\mathrm{x}, \mathrm{t})| 2 \mathrm{dx} \equiv \mathrm{Z}|\psi 0(\mathrm{x})| 2 \mathrm{dx}$ $=1, \mathrm{t} \geq 0$,
and the energy

$$
\begin{align*}
E(\psi(\cdot, t)) & \left.:=\int_{\mathbb{R}^{3}}\left[\psi^{*}-\Delta+m^{2}\right)^{1 / 2} \psi+\left(V_{\mathrm{ext}}(\mathbf{x})+\frac{\lambda}{2|\mathbf{x}|} *|\psi|^{2}\right)|\psi|^{2}\right] \\
& \equiv E\left(\psi_{0}\right), \quad t \geq 0 \tag{dx}
\end{align*}
$$

The well-posedness of the underlying worth issue was widely examined in and references in that. Their outcomes can be abridged as: (I) there exists an all inclusive consistent $\lambda \mathrm{cr}$ (additionally alluded to "as far as possible mass" in material science and with a lower bound $\lambda \mathrm{cr}>4 / \pi$ ) to such an extent that, when $\lambda>-\lambda \mathrm{cr}$, the arrangement is comprehensively well-presented in the vitality space $\mathrm{H} 1 / 2(\mathrm{R} 3)$ gave that $\mathrm{V} \in \mathrm{L} 3(\mathrm{R} 3) \cap$ $L \infty(\mathrm{R} 3)$; (ii) when $\lambda \leq-\lambda \mathrm{cr}$, the arrangement is locally well-presented; and (iii) when $\lambda<-\lambda \mathrm{cr}$, the arrangement will explode in limited time, which shows the "gravitational breakdown" of boson stars when the viable "mass" surpasses the basic esteem $\lambda \mathrm{cr}$. Another issue of interests is the presence and uniqueness of the ground state for which is characterized as the limit of the accompanying no curved minimization issue:

## Find

$\phi_{g} \in S=\left\{\phi \mid \phi \in H^{1 / 2}\left(\mathbb{R}^{3}\right),\|\phi\|^{2}=1\right\}$
such that

$$
\begin{equation*}
\mathrm{Eg}:=\mathrm{E}(\varphi \mathrm{~g})=\min \mathrm{E}(\varphi) . \tag{1.2}
\end{equation*}
$$

If Vext $(x) \equiv 0$, it was shown that the ground state exists iff $-\lambda \mathrm{cr}<\lambda<0$ and any ground state is smooth, decays exponentially when $|\mathrm{x}| \rightarrow \infty$, and is identical to its spherically symmetric rearrangement up to phase and translation. Moreover, it was proven recently in

$$
V_{P}(\mathrm{x}, t)=\frac{1}{4 \pi|\mathbf{x}|} *|\psi|^{2}=\frac{1}{4 \pi} \int_{\mathbb{R}^{3}} \frac{1}{\left|\mathrm{x}-\mathrm{x}^{\prime}\right|}\left|\psi\left(\mathrm{x}^{\prime}, t\right)\right|^{2} \mathrm{~d} \mathbf{x}^{\prime}, \quad \mathrm{x} \in \mathbb{R}^{3}, \quad t \geq 0
$$

Then is re-written as the relativistic Schro"dinger-Poisson (RSP) equation
$i \partial_{t} \psi(\mathbf{x}, t)=\sqrt{-\Delta+m^{2}} \psi+V_{\mathrm{ext}}(\mathbf{x}) \psi+\beta V_{P} \psi, \quad \mathbf{x} \in \mathbb{R}^{3}, \quad t>0$,
$-\Delta V P(x, t)=|\psi| 2, \quad x \in R 3$, xylem VP $(x, t)=0, \quad t \geq 0$.
$|\mid \rightarrow \infty$
With this formulation, the energy functional (1.1 is re-written as

$$
\begin{aligned}
& \left.E(\psi(\cdot, t))=\int_{\mathbb{R}^{3}}\left[\psi^{*}-\Delta+m^{2}\right)^{1 / 2} \psi+\left(V_{\mathrm{ext}}(\mathbf{x})+\frac{\beta}{2} V_{P}\right)|\psi|^{2}\right] \\
= & \left.\left.\int_{\mathbb{R}}\left[\mid-\Delta+m^{2}\right)^{1 / 4} \psi\right|^{2}+\left(V_{\mathrm{ext}}(\mathbf{x})+\frac{\beta}{2}(-\Delta)^{-1}|\psi|^{2}\right)|\psi|^{2}\right] \\
= & \left.\left.\int_{\mathbb{R}^{3}}\left[\mid-\Delta+m^{2}\right)^{1 / 4} \psi\right|^{2}+V_{\mathrm{ext}}(\mathbf{x})|\psi|^{2}+\frac{\beta}{2}\left|\nabla V_{P}\right|^{2}\right]_{\mathrm{dx}} \quad \mathrm{dx}, \quad \mathrm{t} \geq 0
\end{aligned}
$$

dx

So as to plan numerical techniques for processing the ground states, Sine pseudo unearthly technique is connected to defamed it. For processing the elements, again the issue is truncated into a container with homogeneous Dirichlet limit conditions and a period part sine pseudo unearthly technique is connected to discredited it. Specifically, when the potential and beginning information for elements are circularly symmetric, the issue crumples to a semi 1D issue. Like streamlined numerical strategies are planned by utilizing an appropriate difference in factors in the semi 1D issue.

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